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Model calculations of the phase-behaviour of rod-like particles in solution and their application to the isotropic and anisotropic phases of poly(n-alkylisocyanates) in solution

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Aharoni and co-workers¹⁻⁷ have shown that concentrated solutions of certain poly(alkylisocyanates) exhibit lyotropic mesomorphic behaviour. Theoretical considerations of the formation and structure of concentrated solutions of rod-like macromolecules have been given by Flory and co-workers who have shown that anisotropic and isotropic phases may coexist above a minimum concentration of polymer. Recently Flory and co-workers⁸⁻¹⁴ have published an important series of papers in which the earlier work is cited and is extended to the case of polydisperse systems. Specifically, Flory and Frost 10 and Frost and Flory 11 have made calculations for the 'most probable' and Poisson distributions of species, respectively. Their calculations predict that in the biphasic region there will be a partial fractionation of material between the two phases, the higher molecular weight species tending to concentrate in the anisotropic phase. The formation of lyotropic mesophases has been observed by Aharoni and Walsh^{2,3} for certain poly(nalkylisocyanates) in solution, and, further, they showed that under the correct experimental conditions, bi-phasic media involving poly(n-hexylisocyanate) separated into isotropic and anisotropic phases during which there was a partial fractionation of the original polydisperse material in the manner predicted by Flory^{8-11.14}.

During the course of our studies of the dielectric relaxation behaviour of poly(n-hexylisocyanate)-toluene solutions^{15,16} and of a copolymer of n-butyl and n-nonyl isocyanates in toluene solution¹⁷, covering isotropic and bi-phasic ranges, it became apparent that a knowledge, if only semi-quantitative, of the relative amounts of isotropic and anisotropic phases and their compositions would be desirable in order to assist in the interpretations of our data for the reorientations of molecules.

The calculations of Flory and Frost 10,11 were made for distributions having fixed half-width for a given average length of rod. We have extended the work of Flory and Frost 10,11 in order to include the effect of varying the distribution of molecular weight (rod-length) for a given degree of polymerization. This has been done using the theory of Flory and co-workers^{8-11,14} together with a Gaussian distribution of rod-lengths. In this way we may

explore the phase-behaviour as a function of average rodlength, width of distribution and concentration of polymer. The details of the calculations and a complete set of results will be the subject of a further publication 18. The present communication gives a portion of these data and indicates how they may be applied to such systems as the poly (n-alkylisocyanates) in solution.

Theory

The theory of Flory and co-workers⁸⁻¹⁴ is adopted in the present work. We use the same notation as these earlier workers and write the volume-fraction of x-meric species in the system as a whole as v_x^0 and the mean volume-fraction of all polymer (solute) in both phases combined as v_2^0 . The Gaussian distribution for particles consisting of x units leads to the following relation:

$$\frac{v_x^0}{v_2^0} = \frac{x}{\bar{x}} \exp \left[-4 \cdot \ln 2 \cdot \left(\frac{x - x_0}{\Delta_{1/2}} \right)^2 \right]$$
 (1)

where x_0 is the value of x at the maximum of the Gaussian curve, \bar{x} is the mean value of x, $\hat{x} = (\sum x n_x^0 / n_2^0 . n_x^0)$ is the number of x-meric species in the total volume V. The total number of particles in V is given by the normalization condition $n_2^0 = \sum_{1}^{\infty} n_x^0$. The total half-width of the Gaussian distribution, $\Delta_{1/2},$ may be varied for each chosen value of x_0 . Our calculations follow the procedure outlined by Flory and Frost^{10,11} and yield: v_2 and v_2' , being the volume-fractions of solute in the isotropic and anisotropic phases, respectively, and Φ , being the ratio of the volume of the isotropic phase to the total volume, for given values of v_2^0 , x_0 and $\Delta_{1/2}$. The solute concentration c(% w/w) is defined as 100 times the ratio of the weight of solute to the weight of solution, and may be expressed as

$$c = 100 \times \left[1 + \frac{\rho_1}{\rho_2} \left(\frac{1}{v_2^0} - 1 \right) \right]^{-1}$$
 (2)

 ρ_1 and ρ_2 are solvent and solute density respectively. For the present calculations we took $\rho_1/\rho_2 = 0.86/1.00$, to conform with values for poly(n-hexylisocyanate) in toluene³. We calculate $\Phi_A = (1 - \Phi)$, being the volume

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fraction of the anisotropic phase, as a function of c (equation (2)) for $x_0 = 50$ and different values of $\Delta_{1/2}$. This value of x_0 was chosen in order to conform with a rod-like poly(n-hexylisocyanate) at a molecular weight of about 10⁵ (see ref. 19). The correctness of our procedures and calculations were first checked by verification of Table 1 of ref. 10 and Table 1 of ref. 11. All calculations were made using the College Honeywell 6080 main-frame computer.

Figure 1 shows Φ_A for $x_0 = 50$ as a function of c for $\Delta_{1/2}$

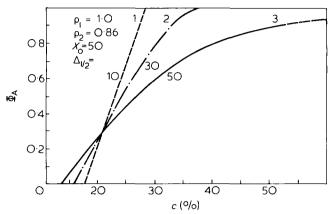


Figure 1 Φ_A against c (% w/w) for $x_0 = 50$ and for different values of $\Delta_{1/2}$. Curves 1, 2 and 3 correspond to $\Delta_{1/2}$ equal to 10, 30 and 50 respectively

equal to 10, 30 and 50. For $\Delta_{1/2} = 10$, the wholly isotropic phase persists up to c = 17.8%. In the range 17.8 < c<28.2% the system is bi-phasic with Φ_A increasing monotonically with c. For c > 28.2% the system is wholly anisotropic. As $\Delta_{1/2}$ is increased two effects are observed: (i) there is a lowering of the concentration at which the biphasic system first appears, (ii) the range of the bi-phasic region is considerably extended. For $\Delta_{1/2} = 30$ the formation of the wholly anisotropic phase is extended up to $c \sim 37\%$.

In addition to the volume fraction of anisotropic phase Φ_A , we obtain the distribution curves for the composition of the isotropic and anisotropic phase for given values of $x_0, \Delta_{1/2}$ and c. From the curves of Figure 1 it is of interest to see the form of these distributions at different values of c. We note that in Figure 1 for c = 21% the values of Φ_A are independent of $\Delta_{1/2}$. Thus we chose to show in Figures 2 the distributions for c below, above and equal to c = 21%. Here f_x and f_x' are the quantities of x-mer in the isotropic and anisotropic phases respectively, as defined by Flory and Frost 10, 11.

$$f_{\mathbf{x}} = \Phi v_{\mathbf{x}} / v_2^0 \tag{3a}$$

$$f_x' = (1 - \Phi)v_x'/v_2^0$$
 (3b)

and

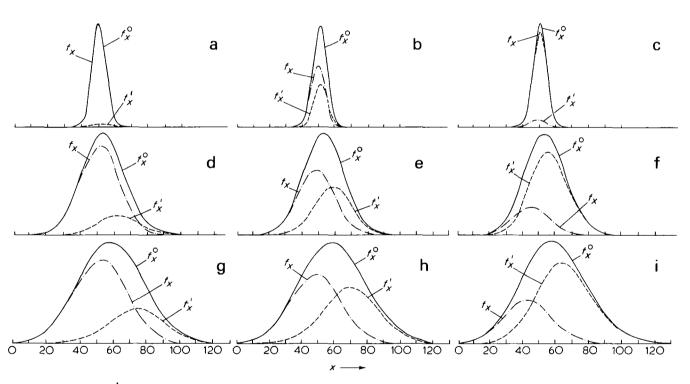


Figure 2 f_X^0 , f_X and f_X' against x for $x_0 = 50$ and different values of $\Delta_{1/2}$ and c

Figure	$\Delta_{1/2}$	c (% w/w)	$^{\Phi}A$
 а	10	18	0.017
b	10	21	0.310
С	10	27	0.890
d	30	18	0.115
е	30	21	0.310
f	30	27	0.640
g	50	18	0.183
h	50	21	0.310
i	50	27	0.520

$$f_{x}^{0} = f_{x} + f_{y}' = v_{x}^{0} / v_{2}^{0}$$
 (4)

Consider first Figure 2e. At this concentration, c = 21%, x_0 = 50 and $\Delta_{1/2}$ = 30. The plots of f_x^0 , f_x and f_x' are each essentially symmetrical about their maximum values. The original mixture partitions into isotropic and anisotropic phases with the high molecular weight species favouring the anisotropic phase and the low molecular weight

Table 1 v_2 , v_2' , $\Phi_A = 1 - \Phi$ and $\Phi_A v_2' / (\Phi v_2)$ for all systems shown in Figure 2

c (% w/w)	$\Delta_1/_2$	v_2	v_2'	Φ	Φ_{A}	$\Phi_{A} v_2'/(\Phi v_2)$
18	10	0.157	0.245	0.983	0.017	0.0270
18	30	0.146	0.250	0.885	0.115	0.222
18	50	0.136	0.260	0.817	0.183	0.429
21	10	0.159	0.248	0.69	0.310	0.700
21	30	0.153	0.262	0.69	0.310	0.769
21	50	0.145	0.279	0.69	0.310	0.864
27	10	0.160	0.252	0.11	0.89	12.74
 27	30	0.164	0.284	0.36	0.637	3.06
 27	50	0.160	0.316	0.48	0.520	2.14

species the isotropic phase. From equation (4)

$$\sum_{x=1}^{\infty} f_x^0 = 1 = \sum_{1}^{\infty} f_x + \sum_{1}^{\infty} f_x'$$

$$= \Phi \frac{v_2}{v_2^0} + \Phi_A \frac{v_2'}{v_2^0}$$
(5)

 v_2 and v_2' are the volume fractions of solute in the isotropic and anisotropic phases respectively. Thus $\Phi_A v_2 / [\Phi v_2]$ is the ratio of the volume of solute in the anisotropic phase to the volume of solute in the isotropic phase. Table 1 lists v_2 , v_2 , Φ , $\Phi_A = 1 - \Phi$, and the above ratio for all systems shown in Figures 2. We note that since the curves of f_x and f'_{x} , for a given system, have essentially the same shape, the ratio of their areas, $\Sigma f_x'/\Sigma f_x = \Phi_A v_2'/[\Phi v_2]$, is given, quite accurately, by the ratio of their peak heights, i.e. $(f_x')_{\text{max}}/(f_x)_{\text{max}}$, so this ratio is readily judged by eye from the figures.

Consider now the sequence—Figures 2b, 2e and 2h, for which Φ_A is constant (see Figure 1). As $\Delta_{1/2}$ is increased the maxima in f_x and f_x' move apart; i.e. as the distribution is broadened the extent of partial fractionation between the phases is increased. Table 1, and the figures, show that the relative amount of polymer in the anisotropic phase increases as $\Delta_{1/2}$ increases although Φ_A is constant.

For the sequence—Figures 2a, 2d and 2g, the system starts as predominantly isotropic phase and ends up as predominantly anisotropic phase, with increasing $\Delta_{1/2}$, but with a steady increase in $\Phi_{A}v_{2}'$ and a corresponding decrease in Φv_2 . As $\Delta_{1/2}$ increases the maxima in f_x and f_x' move apart, as for Figures 2b, 2e and 2h. For the sequence—Figures 2c, 2f and 2i the bi-phasic system starts off as predominantly anisotropic phase but as $\Delta_{1/2}$ increases $\Phi_{\mathsf{A}} v_2'$ steadly decreases with a corresponding increase in Φv_2 . Again the maxima in f_x and f_x' move apart as $\Delta_{1/2}$ is increased. The variation of the amount of anisotropic phase with $\Delta_{1/2}$ for Figures 2c, 2f and 2i contrasts with that for Figures 2a, 2d and 2g. Thus broadening the distribution for a bi-phasic material that is predominantly composed of anisotropic phase (Figures 2c, 2f, 2i) introduces more isotropic phase whereas for a biphasic material that is predominantly composed of isotropic phast (Figures 2a, 2d, 2g) broadening the distribution introduces more anisotropic phase.

We may examine the variations with c at fixed values of $\Delta_{1/2}$. For the sequence—Figures 2a, 2b and 2c, increase in c for a fairly-narrow distribution ($\Delta_{1/2} = 10$) transforms the bi-phasic material from essentially isotropic to essentially anisotropic phase (as expected from Figure 1) with very little partial fractionation with respect to molecular weight. For $\Delta_{1/2} = 30$ Figures 2d, 2e and 2f show that as c is increased $\Phi_A v_2'$ increases rapidly and there is a marked partial fractionation between the phases. Similar behaviour is seen for $\Delta_{1/2} = 50$, i.e. for Figures 2g, 2h and 2i, as c is increased. Note, however, that the maxima in f_x and f'_x move to lower values of x as c is increased for constant $\Delta_{1/2}$. Thus if we have a polymer of a given distribution of molecular weight and form two biphasic materials one of low overall polymer concentration, c_A say, in which the isotropic phase predominates, and one of high concentration, c_B say, in which the anisotropic phase predominates, then the average molecular weight of polymer in the two phases will obey the relations $\langle M_{\text{iso,B}} \rangle < \langle M_{\text{iso,A}} \rangle$, $\langle M_{\text{aniso,B}} \rangle$ $<\langle M_{\rm aniso,A}\rangle$. More comprehensive calculations, including variation of x_0 and their relation to our dielectric data for concentrated solutions of poly(n-hexylisocyanate) will be considered in a future publication 16.

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